# THE DIFFERENTIAL STRUCTURE OF THE BRIESKORN LATTICE 

MATHIAS SCHULZE


#### Abstract

The Brieskorn lattice $H^{\prime \prime}$ of an isolated hypersurface singularity with Milnor number $\mu$ is a free $\mathbb{C}\{\{s\}\}$-module of rank $\mu$ with a differential operator $t=s^{2} \partial_{s}$. Based on the mixed Hodge structure on the cohomology of the Milnor fibre, M. Saito constructed $\mathbb{C}\{\{s\}\}$-bases of $H^{\prime \prime}$ for which the matrix of $t$ has the form $A=A_{0}+A_{1} s$. We describe an algorithm to compute the matrices $A_{0}$ and $A_{1}$. They determine the differential structure of the Brieskorn lattice, the spectral pairs and Hodge numbers, and the complex monodromy of the singularity.


## 1. The Milnor Fibration

Let $f:\left(\mathbb{C}^{n+1}, \underline{0}\right) \longrightarrow(\mathbb{C}, 0)$ be a holomorphic function germ with an isolated critical point and Milnor number $\mu=\operatorname{dim}_{\mathbb{C}} \mathbb{C}\{\underline{x}\} /\langle\underline{\partial}(f)\rangle$ where $\underline{x}=x_{0}, \ldots, x_{n}$ is a complex coordinate system of $\left(\mathbb{C}^{n+1}, 0\right)$ and $\underline{\partial}=\partial_{x_{0}}, \ldots, \partial_{x_{n}}$. By the finite determinacy theorem, we may assume that $f \in \mathbb{C}[\underline{x}]$. By E.J.N Looijenga [7, 2.B], for a good representative $f: X \longrightarrow T$ where $T \subset \mathbb{C}$ is an open disk at the origin, the restriction $f: X^{\prime} \longrightarrow T^{\prime}$ to $T^{\prime}=T \backslash\{0\}$ and $X^{\prime}=X \backslash f^{-1}(0)$ is a $\mathscr{C}^{\infty}$ fibre bundle unique up to diffeomorphism, the Milnor fibration. By J. Milnor [9, 6.5], the general fibre $X_{t}=f^{-1}(t), t \in T^{\prime}$, is homotopy equivalent to a bouquet of $\mu n$-spheres and, in particular, its reduced cohomology is $\widetilde{\mathrm{H}}^{k}\left(X_{t}\right) \cong \delta_{k, n} \mathbb{Z}^{\mu}$ where $\delta$ is the Kronecker symbol. Since $T^{\prime}$ is locally contractible, the $n$-th cohomologies $\mathrm{H}(U)=\mathrm{H}^{n}\left(X_{U}\right)$ of $X_{U}=$ $f^{-1}(U)$ form a locally free $\mathbb{Z}$-sheaf of rank $\mu$ and $\mathrm{H}_{\mathbb{C}}=\mathrm{H} \otimes_{\mathbb{Z}} \mathbb{C}$ is a complex local system of dimension $\mu$. Hence, the sheaf of holomorphic sections $\mathscr{H}=\mathrm{H} \otimes_{\mathbb{Z}} \mathscr{O}_{T^{\prime}}$ of $\mathrm{H}_{\mathbb{C}}$ is a locally free $\mathscr{O}_{T^{\prime}}$-sheaf of rank $\mu$, the cohomology bundle. By P. Deligne [4, 2.23], there is a natural flat connection $\nabla: \mathscr{H} \longrightarrow \mathscr{H} \otimes_{\mathscr{T}_{T^{\prime}}} \Omega_{T^{\prime}}^{1}$ on $\mathscr{H}$ with sheaf of flat sections $\mathrm{H}=\operatorname{ker}(\nabla)$, the Gauss-Manin connection.

## 2. The Monodromy Representation

Let $t$ be a complex coordinate of $T \subset \mathbb{C}, i: T^{\prime} \longrightarrow T$ the canonical inclusion, and $u: T^{\infty} \longrightarrow T^{\prime}$ the universal covering of $T^{\prime}$ defined by $u(\tau)=\exp (2 \pi \mathrm{i} \tau)$ for a complex coordinate $\tau$ of $T^{\infty} \subset \mathbb{C}$. Then the covariant derivative $\nabla_{\partial_{t}}$ of $\nabla$ along $\partial_{t}$ induces a differential operator $\partial_{t}$ on $i_{*} \mathscr{H}$ and the pullback $f^{\infty}: X^{\infty}=X^{\prime} \times_{T^{\prime}} T^{\infty} \longrightarrow T^{\infty}$
is a $\mathscr{C}^{\infty}$ fibre bundle with $X_{\tau}^{\infty}=X_{u(\tau)}$, the (canonical) Milnor fibre. Since $T^{\infty}$ is contractible, the $n$-th cohomologies $H(U)=\mathrm{H}^{n}\left(X_{U}^{\infty}\right)$ of $X_{U}^{\infty}=\left(f^{\infty}\right)^{-1}(U)$ form a free $\mathbb{Z}$-sheaf of rank $\mu$ and $u_{*} H$ is the sheaf of multivalued sections of H . Lifting closed paths in $T^{\prime}$ along sections of H defines the monodromy representation $\pi_{1}\left(T^{\prime}, t\right) \longrightarrow \operatorname{Aut}\left(\mathrm{H}_{t}\right)$ on $\mathrm{H}_{t}$ inducing the monodromy representation $\pi_{1}\left(T^{\prime}\right) \longrightarrow \operatorname{Aut}(H)$ on the cohomology $H$ of the Milnor fibre. The image M of the counterclockwise generator of $\pi_{1}\left(T^{\prime}\right)$ is called the monodromy operator and fulfills $\mathrm{M}(s)(\tau)=s(\tau+1)$ for $s \in H$. The sheaf H is determined by the monodromy representation up to isomorphism. The following well known theorem is due to E. Brieskorn $[2,0.6]$ and others.

Theorem 1 (Monodromy Theorem). The eigenvalues of the monodromy are roots of unity and its Jordan blocks have size at most $(n+1) \times(n+1)$ and size at most $n \times n$ for eigenvalue 1 .

## 3. The Gauss-Manin Connection

Let $\mathrm{M}=\mathrm{M}_{s} \mathrm{M}_{u}$ be the decomposition of M into semisimple part $\mathrm{M}_{s}$ and unipotent part $\mathrm{M}_{u}$ and let $\mathrm{N}=-\frac{\log \mathrm{M}_{u}}{2 \pi \mathrm{i}}$ be the nilpotent part of M. Note that $-2 \pi \mathrm{iN} \in \operatorname{End}_{\mathbb{Q}}\left(H_{\mathbb{Q}}\right)$ where $H_{\mathbb{Q}}=H \otimes_{\mathbb{Z}} \mathbb{Q}$. Let $H_{\mathbb{C}}=\bigoplus_{\lambda} H_{\mathbb{C}}^{\lambda}$ be the decomposition of $H_{\mathbb{C}}=H \otimes_{\mathbb{Z}} \mathbb{C}$ into generalized $\lambda$-eigenspaces $H_{\mathbb{C}}^{\lambda}$ of M and $\mathrm{M}^{\lambda}=\left.\mathrm{M}\right|_{H_{\mathrm{C}}^{\lambda}}$. Note that $H_{\mathrm{Q}}=H_{\mathrm{Q}}^{1} \oplus H_{\mathrm{Q}}^{\neq 1}$ where $H_{\mathbb{Q}}^{1} \otimes_{\mathbb{Q}} \mathbb{C}=H_{\mathbb{C}}^{1}$ and $H_{\mathbb{Q}}^{\neq 1} \otimes_{\mathbb{Q}} \mathbb{C}=\bigoplus_{\lambda \neq 1} H_{\mathbb{C}}^{\lambda}$. Then there is an inclusion

$$
H_{\mathbb{C}}^{\mathrm{e}^{-2 \pi \mathrm{i} \alpha}} \xrightarrow{\psi_{\alpha}}\left(i_{*} \mathscr{H}\right)_{0}
$$

defined by $\psi_{\alpha}(A)=t^{\alpha+\mathrm{N}} A=t^{\alpha} \exp (N \log (t))$ with image $C^{\alpha}=\operatorname{im}\left(\psi_{\alpha}\right)$. In particular, the operators M and N act on $C^{\alpha}$. The following lemma is an immediate consequence of the definition of $\psi_{\alpha}$.

## Lemma 2.

(1) $t \circ \psi_{\alpha}=\psi_{\alpha+1}$ and $\partial_{t} \circ \psi_{\alpha}=\psi_{\alpha-1} \circ(\alpha+\mathrm{N})$.
(2) $t: C^{\alpha} \longrightarrow C^{\alpha+1}$ is bijective and $\partial_{t}: C^{\alpha} \longrightarrow C^{\alpha-1}$ is bijective if $\alpha \neq 0$.
(3) On $C^{\alpha}, t \partial_{t}-\alpha=\mathrm{N}$ and $\exp \left(-2 \pi \mathrm{i} t \partial_{t}\right)=\mathrm{M}^{\mathrm{e}^{-2 \pi \mathrm{i} \alpha}}$.
(4) $C^{\alpha}=\operatorname{ker}\left(t \partial_{t}-\alpha\right)^{n+1}$.

Definition 3. We call $G=\bigoplus_{-1<\alpha \leq 0} \mathbb{C}\{t\}\left[t^{-1}\right] C^{\alpha} \subset\left(i_{*} \mathscr{H}\right)_{0}$ the local Gauss-Manin connection.

The local Gauss-Manin connection is a $\mu$-dimensional $\mathbb{C}\{t\}\left[t^{-1}\right]$-vectorspace and a regular $\mathbb{C}\{t\}\left[\partial_{t}\right]$-module. The generalized $\alpha$-eigenspaces $C^{\alpha}$ of the operator $t \partial_{t}$ define the decreasing filtration on $G$ by free $\mathbb{C}\{t\}$-modules

$$
V^{\alpha}=\bigoplus_{\alpha \leq \beta<\alpha+1} \mathbb{C}\{t\} C^{\beta}, \quad V^{>\alpha}=\bigoplus_{\alpha<\beta \leq \alpha+1} \mathbb{C}\{t\} C^{\beta}
$$

of rank $\mu$, the V-filtration. In contrast to the $\psi_{\alpha}$ and $C^{\alpha}$, the $V^{\alpha}$ are independent of the coordinate $t$. The $C^{\alpha}$ define a splitting

$$
C^{\alpha} \cong V^{\alpha} / V^{>\alpha}=\operatorname{gr}_{V}^{\alpha} G
$$

of the V -filtration and we denote by lead ${ }_{V}$ the leading term with respect to this splitting. The ring $\mathbb{C}\{t\}$ is a free module of rank 1 over the ring

$$
\mathbb{C}\{\{s\}\}=\left\{\sum_{k=0}^{\infty} a_{k} s^{k} \in \mathbb{C} \llbracket s \rrbracket \left\lvert\, \sum_{k=0}^{\infty} \frac{a_{k}}{k!} t^{k} \in \mathbb{C}\{t\}\right.\right\}
$$

where $s=\int_{0}^{1} \mathrm{~d} t$ acts by integration. This fact is generalized by the following lemma [13, 1.3.11].
Lemma 4. The action of $s=\partial_{t}^{-1}$ on $V^{>-1}$ extends to a $\mathbb{C}\{\{s\}\}$-module structure and $V^{>-1}$ is a free $\mathbb{C}\{\{s\}\}$-module of rank $\mu$.

Since $\left[\partial_{t}, t\right]=1,[t, s]=s^{2}$ and hence

$$
t=s^{2} \partial_{s}, \quad \partial_{t} t=s \partial_{s}
$$

We call a free $\mathbb{C}\{\{s\}\}$-submodule of $V^{>-1}$ of rank $\mu$ a $\mathbb{C}\{\{s\}\}$-lattice and call a $t \partial_{t}$-invariant $\mathbb{C}\{\{s\}\}$-lattice saturated. A basis $\underline{e}$ of a $\mathbb{C}\{\{s\}\}$ lattice defines a matrix $A=\sum_{k \geq 0} A_{k} s^{k}$ of $t$ by $t \underline{e}=\underline{e} A$ such that

$$
t \cong A+s^{2} \partial_{s}
$$

is the basis representation of $t$.

## 4. The Brieskorn Lattice

The description of cohomology in terms of holomorphic differential forms by the de Rham isomorphism leads to the definition of the Brieskorn lattice

$$
H^{\prime \prime}=\Omega_{X, 0}^{n+1} / \mathrm{d} f \wedge \mathrm{~d} \Omega_{X, 0}^{n-1} .
$$

By E. Brieskorn [2, 1.5] and M. Sebastiani [15], the Brieskorn lattice is the stalk at 0 of a locally free $\mathscr{O}_{T^{-}}$-sheaf $\mathscr{H}^{\prime \prime}$ of rank $\mu$ with $\left.\mathscr{H}^{\prime \prime}\right|_{T^{\prime}} \cong \mathscr{H}$ and hence $H^{\prime \prime} \subset\left(i_{*} \mathscr{H}\right)_{0}$. The regularity of the Gauss-Manin connection proved by E. Brieskorn $[2,2.2]$ implies that $H^{\prime \prime} \subset G$. B. Malgrange [8, 4.5] improved this result by the following theorem.
Theorem 5. $H^{\prime \prime} \subset V^{-1}$.
By E. Brieskorn [2, 1.5], the Leray residue formula can be used to express the action of $\partial_{t}$ in terms of differential forms by $\partial_{t}[\mathrm{~d} f \wedge \omega]=$ [d $\omega$ ]. In particular, $s H^{\prime \prime} \subset H^{\prime \prime}$ and

$$
H^{\prime \prime} / s H^{\prime \prime} \cong \Omega_{X, 0}^{n+1} / \mathrm{d} f \wedge \Omega_{X, 0}^{n} \cong \mathbb{C}\{\underline{x}\} /\langle\underline{\partial}(f)\rangle .
$$

Since the $V^{>-1}$ is a $\mathbb{C}\{\{s\}\}$-module, theorem 5 implies that $H^{\prime \prime}$ is a free $\mathbb{C}\{\{s\}\}$-module of rank $\mu$ and the action of $s$ can be expressed in terms of differential forms by

$$
s[\mathrm{~d} \omega]=[\mathrm{d} f \wedge \omega] .
$$

For computational purposes, we may restrict our attention to the completion of the Brieskorn lattice. E. Brieskorn [2, 3.4] proved the following theorem.
Theorem 6. The $\mathfrak{m}_{X, 0^{-}}$and $\mathfrak{m}_{T, 0^{-}}$-adic topologies on $H^{\prime \prime}$ coincide.
While the proof of theorem 6 is highly non-trivial, the analogous statement for the $\mathbb{C}\{\{s\}\}$-structure of the Brieskorn lattice is quite elementary [13, 1.5.4].

Proposition 7. The $\mathfrak{m}_{X, 0^{-}}$and $\mathfrak{m}_{\mathbb{C}\{\{s\}\}}$-adic topologies on $H^{\prime \prime}$ coincide.
We call the completion $\widehat{H}^{\prime \prime}$ of $H^{\prime \prime}$ the formal Brieskorn lattice. Since completion is faithfully flat, $\widehat{H}^{\prime \prime}$ is a free $\mathbb{C} \llbracket s \rrbracket$-module of rank $\mu$ with a differential operator $t=s^{2} \partial_{s}$. The equality $[\underline{\partial}(f) \bar{g} \mathrm{~d} \underline{x}]=s[\underline{\partial}(\bar{g}) \mathrm{d} \underline{x}]$ motivates to consider the differential relation $\underline{\partial}(f)-s \underline{\partial}$. It is not difficult to prove that it defines the formal Brieskorn lattice as a quotient of $\mathbb{C} \llbracket s, \underline{x} \rrbracket[13,1.5 .6]$.

## Proposition 8.

$$
\mathbb{C} \llbracket s, \underline{x} \rrbracket \xrightarrow{\pi_{H}} \mathbb{C} \llbracket s, \underline{x} \rrbracket /\langle\underline{\partial}(f)-s \underline{\partial}\rangle \mathbb{C} \llbracket s, \underline{x} \rrbracket \cong \cong_{\mathbb{C} \llbracket s \rrbracket} \widehat{H}^{\prime \prime} .
$$

Proposition 8 is the starting point for an algorithmic approach to the local Gauss-Manin connection. Let $<_{\underline{x}}$ be a local degree ordering on $\mathbb{C} \llbracket \underline{x} \rrbracket$ such that $\operatorname{deg}(\underline{x})<\underline{0}$ and $\operatorname{deg}(\underline{\partial})=-\operatorname{deg}(\underline{x})>\underline{0}$. One can compute a polynomial standard basis $\underline{g}$ of the Jacobian ideal $\langle\underline{\partial}(f)\rangle$ and a polynomial transformation matrix $B=\left(\bar{b}^{j}\right)^{j}$ such that $\underline{g}=\underline{\partial}(f) B$. By Nakayama's lemma, $\underline{m}=\left(\underline{x}^{\underline{\beta}}\right)_{\underline{\underline{\beta}} \underline{\underline{\beta}} \notin(\operatorname{lead}(\underline{g})\rangle}$ represents a $\mathbb{C} \llbracket s \rrbracket$-basis $[\underline{m}]$ of $\widehat{H}^{\prime \prime}$. Let $<_{s}$ be the local degree ordering on $\mathbb{C} \llbracket s \rrbracket$ and let $<=\left(<_{s},<_{\underline{x}}\right)$ be the block ordering of $<_{s}$ and $<_{\underline{x}}$ on $\mathbb{C} \llbracket s, \underline{x} \rrbracket$.

## Definition 9.

(1) $\underline{h}=\left(\left(g_{j}-s \underline{\partial} \bar{b}^{j}\right) \underline{x}^{\underline{\beta}}\right)_{j, \underline{\beta}}$.
(2) $\operatorname{deg}(s)=\min \operatorname{deg}(\underline{m})+2 \min \operatorname{deg}(\underline{x})<0$.
(3) $N=\left(N_{K}\right)_{K \geq 0}$ with $N_{K}=K \operatorname{deg}(s)-2 \min \operatorname{deg}(\underline{x})$.
(4) $V=\left(V_{K}\right)_{K \geq 0}$ with $V_{K}=\left\{p \in \mathbb{C} \llbracket s, \underline{x} \rrbracket \mid \operatorname{deg}(p)<N_{K}\right\}+\langle s\rangle^{K} \subset$ $\mathbb{C} \llbracket s, \underline{x} \rrbracket$.
Since $\widehat{H}^{\prime \prime}$ is a free $\mathbb{C} \llbracket s \rrbracket$-module, $\underline{h}$ is a standard basis of the $\mathbb{C} \llbracket s \rrbracket$ module $\langle\underline{\partial}(f)-s \underline{\partial}\rangle \mathbb{C} \llbracket s, \underline{x} \rrbracket$. The following lemma is technical but not very deep and can be generalized to formal differential deformations [13, 2.2.10].

Lemma 10. $V=\left(V_{K}\right)_{K \geq 0}$ is a basis of the $\langle s, \underline{x}\rangle$-adic topology of $\mathbb{C} \llbracket s, \underline{x} \rrbracket$ with $\pi_{H}\left(V_{K}\right)=\langle s\rangle^{K} \widehat{H}^{\prime \prime}$. If $s^{\alpha} \operatorname{lead}\left(h_{j, \underline{\beta}}\right) \in V_{K}$ then $s^{\alpha} h_{j, \underline{\beta}} \in V_{K}$.

Lemma 10 leads to a normal form algorithm for the Brieskorn lattice [13, 2.2.12]. It computes a normal form with respect to $\underline{h}$ and hence
the [ $\underline{m}$ ]-basis representation in $H^{\prime \prime}$. The normal form computation up to a given degree can be continued up to any higher degree without additional computational effort. The normal form algorithm for the Brieskorn lattice is a special case of a modification of Buchberger's normal form algorithm [3] for power series rings where termination is replaced by adic convergence [13, 2.1.19].

## 5. Mixed Hodge Structure

By lemma 2, there is a $\mathbb{C}$-isomorphism

$$
H_{\mathrm{C}}=\bigoplus_{-1<\alpha \leq 0} H_{\mathbb{C}}^{\mathrm{e}^{-2 \pi \mathrm{i} \alpha}} \xrightarrow{\psi} \bigoplus_{-1<\alpha \leq 0} C^{\alpha} \cong V^{>-1} / s V^{>-1}
$$

defined by $\psi=\bigoplus_{-1<\alpha \leq 0} \psi_{\alpha}$ and the monodromy M on $H_{\mathbb{C}}$ corresponds to $\exp \left(-2 \pi \mathrm{i} t \partial_{t}\right)$ on $\bigoplus_{-1<\alpha \leq 0} C^{\alpha}$.

The Hodge filtration $F=\left(F_{k}\right)_{k \in \mathbb{Z}}$ on $V^{>-1}$ defined by J. Scherk and J.H.M. Steenbrink [14] is the increasing filtration by the free $\mathbb{C}\{\{s\}\}$ modules

$$
F_{k}=F^{n-k}=\left(s^{-k} H^{\prime \prime}\right) \cap V^{>-1}
$$

of rank $\mu$. Via the splitting $C^{\alpha} \cong \operatorname{gr}_{V}^{\alpha} V^{>-1}$, the Hodge filtration induces an increasing Hodge filtration $F C^{\alpha}$ by $\mathbb{C}$-vectorspaces on $C^{\alpha}$ and, via $\psi$, on $H_{\mathbb{C}}$. The nilpotent operator $-2 \pi \mathrm{iN} \in \operatorname{End}_{\mathbb{Q}}\left(H_{\mathbb{Q}}\right)$ defines an increasing weight filtration $W=\left(W_{k}\right)_{k \in \mathbb{Z}}$ centered at $n$ resp. $n+1$ on $H_{\mathrm{Q}}^{\neq 1}$ resp. $H_{\mathrm{Q}}^{1}$.
Theorem 11. The weight filtration $W$ on $H_{\mathrm{Q}}$ and the Hodge filtration $F$ on $H_{\mathbb{C}}$ define a mixed Hodge structure on the cohomology $H$ of the Milnor fibre and the operator N is a morphism of mixed Hodge structures of type $(-1,-1)$.

The mixed Hodge structure on the cohomology of the Milnor fibre was discovered by J.H.M. Steenbrink [16] and described in terms of the Brieskorn lattice by A.N. Varchenko [17].

The nilpotent operator $N$ on $C^{\alpha}$ defines an increasing weight filtration $W=\left(W_{k}\right)_{k \in \mathbb{Z}}$ centered at $n$ on $C^{\alpha}$. By definition N commutes with $\psi_{\alpha}$ and hence

$$
\psi_{\alpha}\left(W H_{\mathbb{C}}^{\mathrm{e}^{-2 \pi \mathrm{i} \alpha}}\right)= \begin{cases}W C^{\alpha}, & \alpha \notin \mathbb{Z} \\ W[-1] C^{\alpha}, & \alpha \in \mathbb{Z}\end{cases}
$$

The weight filtration $W=\bigoplus_{-1<\alpha<0} \mathbb{C}\{\{s\}\} W C^{\alpha}$ on $V^{>-1}$ by free $\mathbb{C}\{\{s\}\}$-modules induces $W C^{\alpha}$ via the splitting $C^{\alpha} \cong \operatorname{gr}_{V}^{\alpha} V^{>-1}$.

The spectral pairs are those pairs $(\alpha, l) \in \mathbb{Q} \times \mathbb{Z}$ with positive multiplicity

$$
d_{l}^{\alpha}=\operatorname{dim}_{\mathbb{C}} \operatorname{gr}_{l}^{W} \operatorname{gr}_{V}^{\alpha} \operatorname{gr}_{0}^{F} V^{>-1}
$$

Via the isomorphism $\psi$, they correspond to the Hodge numbers

$$
h_{\lambda}^{p, l-p}=\operatorname{dim}_{\mathbb{C}} \operatorname{gr}_{F}^{p} \operatorname{gr}_{l}^{W} H_{\mathbb{C}}^{\lambda}
$$

by $d_{l}^{\alpha+p}=h_{\mathrm{e}^{-2 \pi \mathrm{i} \alpha}}^{n-p, l-n+p}$ for $-1<\alpha<0$ and $d_{l}^{p}=h_{1}^{n-p, l+1-n+p}$ and inherit the symmetry properties

$$
d_{l}^{\alpha}=d_{l}^{2 n-l-1-\alpha}, \quad d_{l}^{\alpha}=d_{2 n-l}^{\alpha-n+l}, \quad d_{l}^{\alpha}=d_{2 n-l}^{n-1-\alpha}
$$

from the mixed Hodge structure. The spectral numbers are those numbers $\alpha \in \mathbb{Q}$ with positive multiplicity

$$
d^{\alpha}=\operatorname{dim}_{\mathbb{C}} \operatorname{gr}_{V}^{\alpha} \operatorname{gr}_{0}^{F} V^{>-1}=\sum_{l \in \mathbb{Z}} d_{l}^{\alpha}
$$

and have the symmetry property $d^{\alpha}=d^{n-1-\alpha}$.

## 6. M. Saito's Basis

By P. Deligne [5, 1.2.8], a morphism of mixed Hodge structures is strict for the Hodge filtration. In particular, by theorem 11, N is strict for the Hodge filtration on $H_{\mathbb{C}}$ and on $\mathrm{gr}_{V} V^{>-1}$. Hence, there is a direct sum decomposition $F_{k} C^{\alpha}=\bigoplus_{j \leq k} C^{\alpha, j}$ such that $\mathrm{N}\left(C^{\alpha, k}\right) \subset C^{\alpha, k+1}$, and $s C^{\alpha, k} \subset C^{\alpha+1, k-1}$. By definition of the Hodge filtration,

$$
\operatorname{lead}_{V}\left(H^{\prime \prime}\right)=\sum_{\alpha \in \mathbb{Q}} \sum_{k \leq 0} \mathbb{C}\{\{s\}\} C^{\alpha, k}=\bigoplus_{\alpha \in \mathbb{Q}} \mathbb{C}\{\{s\}\} G^{\alpha}
$$

where $G^{\alpha}=C^{\alpha, 0}$. Let $<_{\mathbb{Q} \times \mathbb{Z}}=\left(>_{\mathbb{Q}},>_{\mathbb{Z}}\right)$ be the block ordering of $>_{\mathbb{Q}}$ and $>_{\mathbb{Z}}$ on the index set $\mathbb{Q} \times \mathbb{Z}$. Then the Hodge filtration defines a refinement of the V -filtration on $V^{>-1}$ by free $\mathbb{C}\{\{s\}\}$-modules $V^{\alpha, k}=F_{k} C^{\alpha} \oplus V^{>\alpha}$ of rank $\mu$ and the $C^{\alpha, k}$ define a splitting of this refined filtration compatible with $s$. We call the refinement the Hodge refinement and the splitting a Hodge splitting. The following lemma follows essentially from the fact that $\mathbb{C}\{\{s\}\}$ is a discrete valuation ring [13, 1.10.5,1.10.10].
Lemma 12. Let $H$ be a $\mathbb{C}\{\{s\}\}$-lattice and $C^{\alpha, k}$ a splitting of a refinement of the $V$-filtration compatible with $s$. Then a minimal standard basis of $H$ is a $\mathbb{C}\{\{s\}\}$-basis and there is a reduced minimal standard basis of $H$.

In particular, there is a reduced minimal standard basis of $H^{\prime \prime}$ for a Hodge splitting. The following proposition follows essentially from lemma 2.3 [13, 1.10.12].
Proposition 13. Let $\underline{h}$ be a reduced minimal standard basis of $H^{\prime \prime}$ for a Hodge splitting. Then the $\underline{\text { h-matrix } A}$ of t has degree 1. In particular,

$$
\left(H^{\prime \prime}, t\right) \stackrel{\underline{h}}{\leftrightarrows}\left(\mathbb{C}\{\{s\}\}^{\mu}, A_{0}+A_{1} s+s^{2} \partial_{s}\right)
$$

is an isomorphism. Moreover, $A_{1}$ is semisimple with eigenvalues the spectral numbers of $f$ added by 1 and $\operatorname{gr}_{V}\left(A_{0}\right)$ can be identified with N .

Note that the matrices $A_{0}$ and $A_{1}$ in proposition 13 determine the differential structure of the Brieskorn lattice. M. Saito [10] first constructed a $\mathbb{C}\{\{s\}\}$-basis of $H^{\prime \prime}$ as in proposition 13 without calling it a reduced minimal standard basis.

## 7. The Algorithm

We describe an algorithm to compute $A_{0}$ and $A_{1}$ as in proposition 13 [13]. This algorithm can be simplified to compute the complex monodromy, the spectral numbers, or the spectral pairs only [13].

The normal form algorithm for the Brieskorn lattice in section 4 computes the $[\underline{m}]$-matrix $A=\sum_{k \geq 0} A_{k} s^{k}$ of $t$ defined by $t[\underline{m}]=[f \underline{\mathrm{~m}}]=$ $[\underline{m}] A$ up to any degree. We identify the columns of a matrix $H$ with the generators of a submodule $\langle H\rangle \subset \mathbb{C} \llbracket s \rrbracket^{\mu}$ and denote by $E$ the unit matrix. Then $\langle E\rangle$ is the $[\underline{m}]$-basis representation of $\widehat{H}^{\prime \prime}$. Hence, the following two statements hold for $\underline{h}=[\underline{m}]$ with $\kappa=0$ and $H=E$.
( $H_{\underline{h}}$ ) One can compute $\kappa \geq 0$ and a $\mu \times \mu$-matrix $H$ with coefficients in $\mathbb{C}[s]$ of degree at most $\kappa$ such that $\langle H\rangle$ is the $\underline{h}$-basis representation of $\widehat{H}^{\prime \prime}$ and $s^{\kappa}\langle E\rangle \subset\langle H\rangle$.
$\left(A_{\underline{h}}\right)$ One can compute the $\underline{h}$-matrix $A$ of $t$ up to any degree.
Step by step, we improve the $\mathbb{C} \llbracket s \rrbracket$-basis $\underline{h}$ and show that $\left(H_{\underline{h}}\right)$ and $\left(A_{\underline{\underline{h}}}\right)$ hold. After the last step, $A_{0}$ and $A_{1}$ as in proposition 13 can be computed by a basis transformation of $A$ to a reduced minimal standard basis of $\langle H\rangle$ up to a certain degree bound.

We call the canonical projection jet ${ }_{k}: \mathbb{C} \llbracket s \rrbracket \longrightarrow \bigoplus_{j=0}^{k} \mathbb{C} s^{j}$ the $k$ jet. Let the monomial ordering on $\mathbb{C} \llbracket s \rrbracket^{\mu}=\mathbb{C} \llbracket s \rrbracket \otimes_{\mathbb{C}} \mathbb{C}^{\mu}$ be the block ordering $<=\left(<_{s},>_{\mu}\right)$ of the local degree ordering $<_{s}$ on $\mathbb{C} \llbracket s \rrbracket$ and the inverse ordering $>_{\mu}$ on the indices of the basis elements of $\mathbb{C}^{\mu}$.
7.1. The Saturation of $H^{\prime \prime}$. In this step, we show that $\left(H_{\underline{\underline{h}}}\right)$ and $\left(A_{\underline{\underline{h}}}\right)$ hold for a $\mathbb{C} \llbracket s \rrbracket$-basis $\underline{h}$ of a saturated $\mathbb{C} \llbracket s \rrbracket$-lattice.

The increasing sequence of $\mathbb{C} \llbracket s \rrbracket$-lattices defined by

$$
\widehat{H}_{0}^{\prime \prime}=\widehat{H}^{\prime \prime}, \quad \widehat{H}_{k+1}^{\prime \prime}=s \widehat{H}_{k}^{\prime \prime}+t \widehat{H}_{k}^{\prime \prime} \subset \widehat{H}^{\prime \prime}
$$

is stationary since $\widehat{H}^{\prime \prime}$ is noetherian. Hence, the saturation $\widehat{H}_{\infty}^{\prime \prime}=$ $\bigcup_{k \geq 0} \widehat{H}_{k}^{\prime \prime}$ of $\widehat{H}^{\prime \prime}$ is a saturated $\mathbb{C} \llbracket s \rrbracket$-lattice. The $[\underline{m}]$-basis representation $\left\langle H_{k}\right\rangle$ of $\widehat{H}_{k}^{\prime \prime}$ can be computed by

$$
H_{0}=Q_{-1}=E, \quad Q_{k}=\left(\operatorname{jet}_{k}(A)+s^{2} \partial_{s}\right) Q_{k-1}, \quad H_{k+1}=\left(s H_{k} \mid Q_{k}\right)
$$

We successively compute the $H_{k}$ and check in each step if $\left\langle Q_{k}\right\rangle \subset\left\langle H_{k}\right\rangle$ by a standard basis and normal form computation. If $\left\langle Q_{k}\right\rangle \subset\left\langle H_{k}\right\rangle$ then we stop the computation and set $\kappa=k$ and $H_{\infty}=H_{\kappa}$. Then $\left\langle H_{\infty}\right\rangle$ is the $[\underline{m}]$-basis representation of $\widehat{H}_{\infty}^{\prime \prime}$. We replace $H_{\infty}$ by a minimal standard basis of $\left\langle H_{\infty}\right\rangle$. Then $\underline{h}=s^{-\kappa} \underline{h} H_{\infty}$ is a $\left.\mathbb{C} \llbracket s\right]$-basis of a saturated $\mathbb{C} \llbracket s \rrbracket$-lattice. By a normal form computation with respect to
$H_{\infty}$ up to degree $\kappa$, we compute the $\underline{h}$-basis representation $\left\langle H_{\infty}^{-1} s^{\kappa} E\right\rangle=$ $\left\langle\mathrm{jet}_{\kappa}\left(H_{\infty}^{-1} s^{\kappa} E\right)\right\rangle$ of $\widehat{H}^{\prime \prime}$. Since $\left\langle H_{\infty}\right\rangle \subset\langle E\rangle, s^{\kappa}\langle E\rangle \subset\left\langle H_{\infty}^{-1} s^{\kappa} E\right\rangle$. By a normal form computation with respect to $H_{\infty}$ up to degree $\kappa+k$, one can compute the $k$-jet
$\operatorname{jet}_{k}\left(H_{\infty}^{-1}\left(A-\kappa s E+s^{2} \partial_{s}\right) H_{\infty}\right)=\operatorname{jet}_{k}\left(H_{\infty}^{-1}\left(\operatorname{jet}_{\kappa+k}(A-\kappa s E)+s^{2} \partial_{s}\right) H_{\infty}\right)$
of the $\underline{h}$-matrix of $t$ for any $k \geq 0$.
7.2. The V-Filtration. In this step, we show that $\left(H_{\underline{v}}\right)$ and $\left(A_{\underline{v}}\right)$ hold for a $<_{\mathbb{Q}}$-increasingly ordered $\mathbb{C} \llbracket s \rrbracket$-basis $\underline{v}$ of a $\widehat{V}^{\alpha}$ compatible with the direct sum decomposition $\widehat{V}^{\alpha} / s \widehat{V}^{\alpha} \cong \bigoplus_{\alpha \leq \beta<\alpha+1} C^{\beta}$.

Since $\underline{h}$ is a $\mathbb{C} \llbracket s \rrbracket$-basis of a saturated $\mathbb{C}\{\{s\}\}$-lattice, $A_{0}=0$ and, by theorem 1, the eigenvalues of $A_{1}$ are rational. In order to compute the eigenvalues of $A_{1}$, we transform $A_{1}$ to Hessenberg form and factorize the characteristic polynomials of its blocks. Then we compute a constant $\mathbb{C} \llbracket s \rrbracket$-basis transformation such that $A_{1}=\operatorname{diag}\left(\alpha_{1}, \ldots, \alpha_{\mu}\right)+N$ with $\alpha_{1} \leq \cdots \leq \alpha_{\mu}$ where $\operatorname{diag}\left(\alpha_{1}, \ldots, \alpha_{\mu}\right)$ denotes the diagonal matrix with entries $\alpha_{1}, \ldots, \alpha_{\mu}$. If $\alpha_{\mu}-\alpha_{1}<1$ then $\underline{v}=\underline{h}$ is a $<_{\mathbb{Q}}$-increasingly ordered $\mathbb{C} \llbracket s \rrbracket$-basis $\underline{v}$ of a $\widehat{V}^{\alpha}$ compatible with the direct sum decomposition $\widehat{V}^{\alpha} / s \widehat{V}^{\alpha} \cong \bigoplus_{\alpha \leq \beta<\alpha+1} C^{\beta}$. If $\alpha_{\mu}-\alpha_{1} \geq 1$ then we proceed as follows. Let

$$
A=\left(\begin{array}{ll}
A^{1,1} & A^{1,2} \\
A^{2,1} & A^{2,2}
\end{array}\right)
$$

such that $A_{0}=0, A_{1}^{1,2}=0, A_{1}^{2,1}=0$, and the eigenvalues of $A_{1}^{1,1}$ are the eigenvalues $\alpha$ of $A_{1}$ with $\alpha<\alpha_{1}+1$. Then the $\mathbb{C} \llbracket s \rrbracket\left[s^{-1}\right]$-basis transformation

$$
H \mapsto\left(\begin{array}{cc}
\frac{1}{s} & 0 \\
0 & 1
\end{array}\right) H, \quad A \mapsto\left(\begin{array}{cc}
\frac{1}{s} & 0 \\
0 & 1
\end{array}\right)\left(A+s^{2} \partial_{s}\right)\left(\begin{array}{ll}
s & 0 \\
0 & 1
\end{array}\right)=\left(\begin{array}{cc}
A^{1,1}+s & \frac{1}{s} A^{1,2} \\
s A^{2,1} & A^{2,2}
\end{array}\right)
$$

decreases $\alpha_{\mu}-\alpha_{1}$ and the degree up to which $A$ is computed by 1 and increases $\kappa$ by 1 . After at most $n$ such transformations, $\alpha_{\mu}-\alpha_{1}<1$.
7.3. The Canonical V-Splitting. In this step, we show that $\left(H_{\underline{c}}\right)$ and $\left(A_{\underline{c}}\right)$ hold for a $<_{\mathbb{Q}}$-increasingly ordered $\mathbb{C}$-basis $\underline{c}$ of a direct sum $\bigoplus_{\alpha \leq \beta<\alpha+1} C^{\beta}$ compatible with the direct sum.

Let $\underline{c}$ be the image of $[\underline{v}]$ under the splitting $\widehat{V}^{\alpha} / s \widehat{V}^{\alpha} \cong \bigoplus_{\alpha \leq \beta<\alpha+1} C^{\beta}$. By Nakayama's lemma, $\underline{c}$ is a $\mathbb{C}$-basis of $\bigoplus_{\alpha \leq \beta<\alpha+1} C^{\beta}$ compatible with the direct sum. The eigenvalues of the commutator $\left[\cdot, A_{1}\right] \in \operatorname{End}_{\mathbb{C}}\left(\mathbb{C}^{\mu^{2}}\right)$ are the differences of the eigenvalues of $A_{1}$. Since $\alpha_{\mu}-\alpha_{1}<1$, $\left[\cdot, A_{1}\right]-k \in \mathrm{GL}_{\mu^{2}}(\mathbb{C})$ for $k \geq 1$. Let $U=\sum_{j=0}^{\infty} U_{j} s^{j}$ be the $\mathbb{C} \llbracket s \rrbracket-$ basis transformation defined by $\underline{c}=\underline{v} U$. Then $U_{0}=E$ and $U A_{1} s=$ $\left(A+s^{2} \partial_{s}\right) U$ or equivalently

$$
U_{k}=\left(\left[\cdot, A_{1}\right]-k\right)^{-1} \sum_{j=0}^{k-1} A_{k-j+1} U_{j}
$$

for $k \geq 1$ and hence one can compute $U$ up to any degree. Since $U_{0}=E$ and $\kappa \geq 0$, jet ${ }_{\kappa}(U)$ is a minimal standard basis of $\langle E\rangle$. By a normal form computation with respect to $U$ up to degree $\kappa$, we compute the $\underline{c}$-basis representation $\left\langle U^{-1} H\right\rangle=\left\langle\operatorname{jet}_{\kappa}\left(\operatorname{jet}_{\kappa}(U)^{-1} H\right)\right\rangle$ of $\widehat{H}^{\prime \prime}$ and $A_{1} s$ is the $\underline{c}$-matrix of $t$.

### 7.4. A Hodge Splitting.

In this step, we show that $\left(H_{\underline{f}}\right)$ and $\left(A_{\underline{f}}\right)$ hold for a $<_{\mathbb{Q} \times \mathbb{Z}}$-decreasingly ordered $\mathbb{C}$-basis $\underline{f}$ of a direct sum $\bigoplus_{\alpha \leq \beta<\alpha+1} \bigoplus_{k \in \mathbb{Z}} C^{\beta, k}$ compatible with the direct sum and that one can compute $A_{0}$ and $A_{1}$ as in proposition 13 .

We compute a standard basis of $H$ up to degree $\kappa$ in order to compute the $\underline{c}$-basis representation of the Hodge filtration $F$. The nilpotent part of $A_{1}$ is the $\underline{c}$-basis representation of N . By computing images and quotients of $\mathbb{C}$-vectorspaces, we compute the $\underline{c}$-basis representation of a Hodge splitting $F_{k} C^{\beta}=\bigoplus_{j \leq k} C^{\beta, j}$. Then we compute a constant $\mathbb{C} \llbracket s \rrbracket$ basis transformation $\underline{f}=\underline{c} U$ such that $\underline{f}$ is a $<_{\mathbb{Q} \times \mathbb{Z}}$-decreasingly ordered C-basis of the direct sum $\bigoplus_{\alpha \leq \beta<\alpha+1} \bigoplus_{k \in \mathbb{Z}} C^{\beta, k}$ compatible with the direct sum.

We replace $H$ by a reduced minimal standard basis of $\langle H\rangle$ up to degree $\kappa+1$. By a normal form computation with respect to $H$ up to degree $\kappa+1$, we compute the 1 -jet
$\operatorname{jet}_{1}\left(H^{-1}\left(A+s^{2} \partial_{s}\right) H\right)=\operatorname{jet}_{1}\left(\operatorname{jet}_{\kappa+1}(H)^{-1}\left(\operatorname{jet}_{\kappa+1}(A)+s^{2} \partial_{s}\right) \operatorname{jet}_{\kappa+1}(H)\right)$
of the $\underline{c} H$-matrix $A$ of $t$ in order to compute $A_{0}$ and $A_{1}$ as in proposition 10.

## 8. An Example

The algorithm in section 7 is implemented in the computer algebra system Singular [6] in the procedure tmatrix in the library gaussman.lib [12]. In an example Singular session, we compute the differential structure of the Brieskorn lattice of the singularity of type $T_{2,5,5}$ defined by the polynomial $f=x^{2} y^{2}+x^{5}+y^{5}$.

First, we load the Singular library gaussman.lib:

```
> LIB "gaussman.lib";
```

Then, we define the local ring $R=\mathbb{Q}[x, y]_{\langle x, y\rangle}$ with the local degree ordering ds as monomial ordering and the polynomial $f=x^{2} y^{2}+x^{5}+$ $y^{5} \in R$ :

```
> ring R=0,(x,y),ds;
> poly f=x2y2+x5+y5;
```

Finally, we compute $A_{0}$ and $A_{1}$ as in proposition 10:

```
> list A=tmatrix(f);
```

The result is the list $\mathrm{A}=\mathrm{A}[1], \mathrm{A}[2]$ such that $\mathrm{A}[\mathrm{i}+1]=A_{i}$ and
$A_{0}=\left(\begin{array}{cccc}0 & 0 & \cdots & 0 \\ \vdots & \vdots & & \vdots \\ 0 & 0 & \cdots & 0 \\ 1 & 0 & \cdots & 0\end{array}\right), \quad A_{1}=\operatorname{diag}\left(\frac{1}{2}, \frac{7}{10}, \frac{7}{10}, \frac{9}{10}, \frac{9}{10}, 1, \frac{11}{10}, \frac{11}{10}, \frac{13}{10}, \frac{13}{10}, \frac{3}{2}\right)$.
By proposition $10,\left(H^{\prime \prime}, t\right) \cong\left(\mathbb{C}\{\{s\}\}^{\mu}, A_{0}+s A_{1}+s^{2} \partial_{s}\right)$ and the spectral pairs are $\left(-\frac{1}{2}, 2\right),\left(-\frac{3}{10}, 1\right)^{2},\left(-\frac{1}{10}, 1\right)^{2},(0,1),\left(\frac{1}{10}, 1\right)^{2},\left(\frac{3}{10}, 1\right)^{2}$, $\left(\frac{1}{2}, 0\right)$.

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M. Schulze, Department of Mathematics, D-67653 Kaiserslautern E-mail address: mschulze@mathematik.uni-kl.de

